

Physics 661

Particle Physics Phenomenology

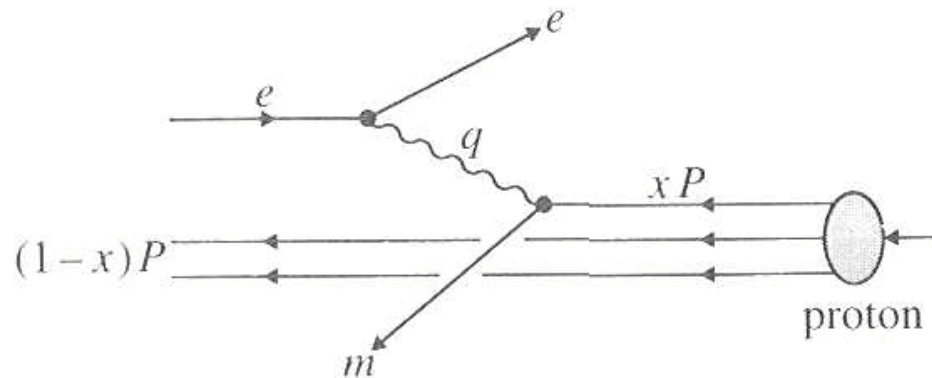
December 2, 2003

Lepton and Quark Scattering

- $e^+ e^- \rightarrow \mu^+ \mu^-$
- $e^+ e^-$ annihilation to hadrons ($e^+ e^- \rightarrow Q\bar{Q}$)
- Electron-muon scattering, $e^- \mu^+ \rightarrow e^- \mu^+$
- Neutrino-electron scattering, $\nu_e e^- \rightarrow \nu_e e^-$
- Elastic lepton-nucleon scattering
- Deep inelastic scattering and partons
- Deep inelastic scattering and quarks
 - Electron-nucleon scattering
 - Neutrino-nucleon scattering
- Quark distributions within the nucleon
- Sum rules

The "discovery" of quarks

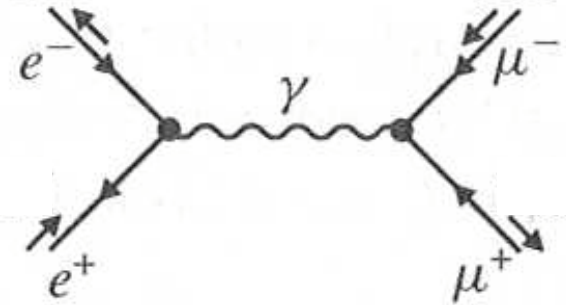
- deep inelastic lepton-nucleon scattering revealed dynamical understanding of quark substructure



- leptonproduction of hadrons could be interpreted as elastic scattering of the lepton by a pointlike constituent of the nucleon, the quark
- theory of scattering of two spin-1/2, pointlike particles required

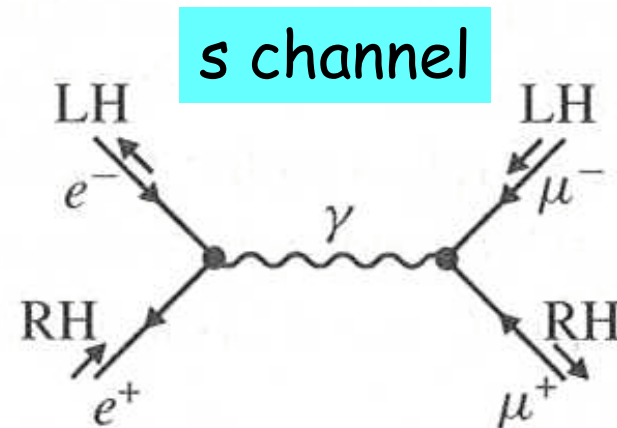
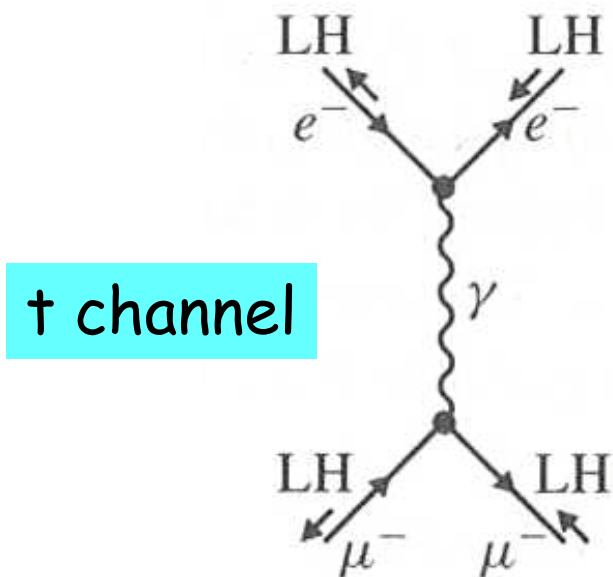
$$e^+ e^- \rightarrow \mu^+ \mu^-$$

- Dominated by single-photon exchange
- $M_{if} = e^2/q^2 = 4\pi\alpha/q^2$
- $\frac{d\sigma}{d\Omega} = \frac{\alpha^2}{4s} f(\theta)$
- What about spins?
 - The conservation of helicity at high energy for the EM interaction means only LR and RL states will interact

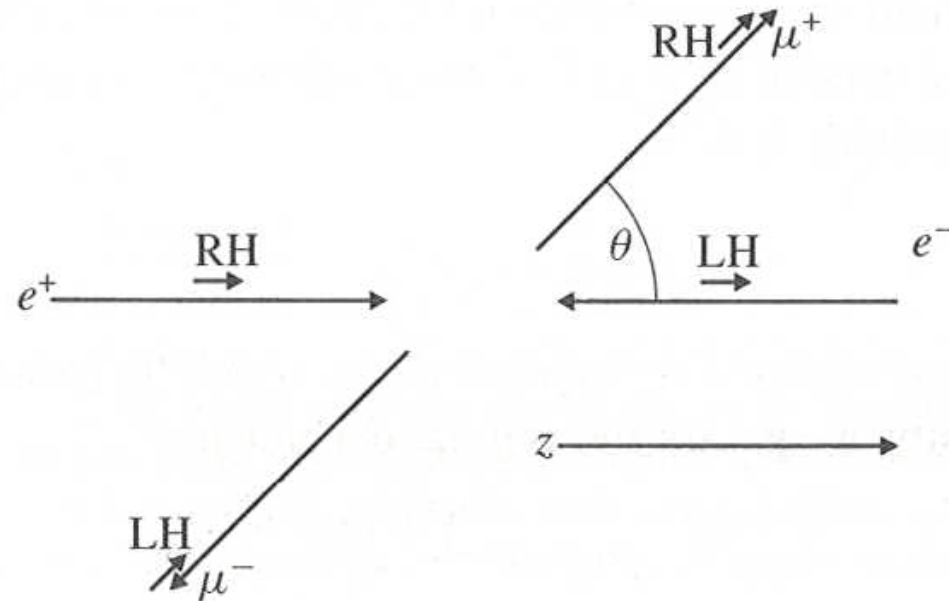


$$e^+ e^- \rightarrow \mu^+ \mu^-$$

- Conservation of helicity
 - consider crossed diagrams
 - (recall, helicity is conserved in vector int. at high energy)



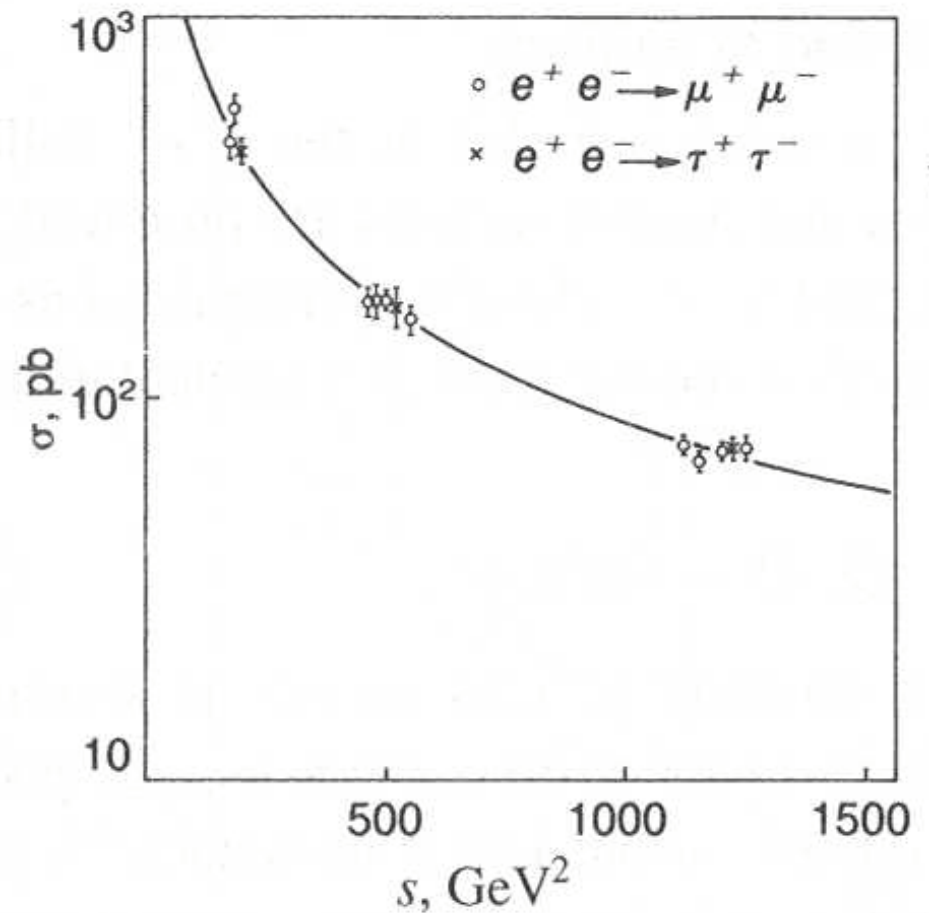
$$e^+ e^- \rightarrow \mu^+ \mu^-$$



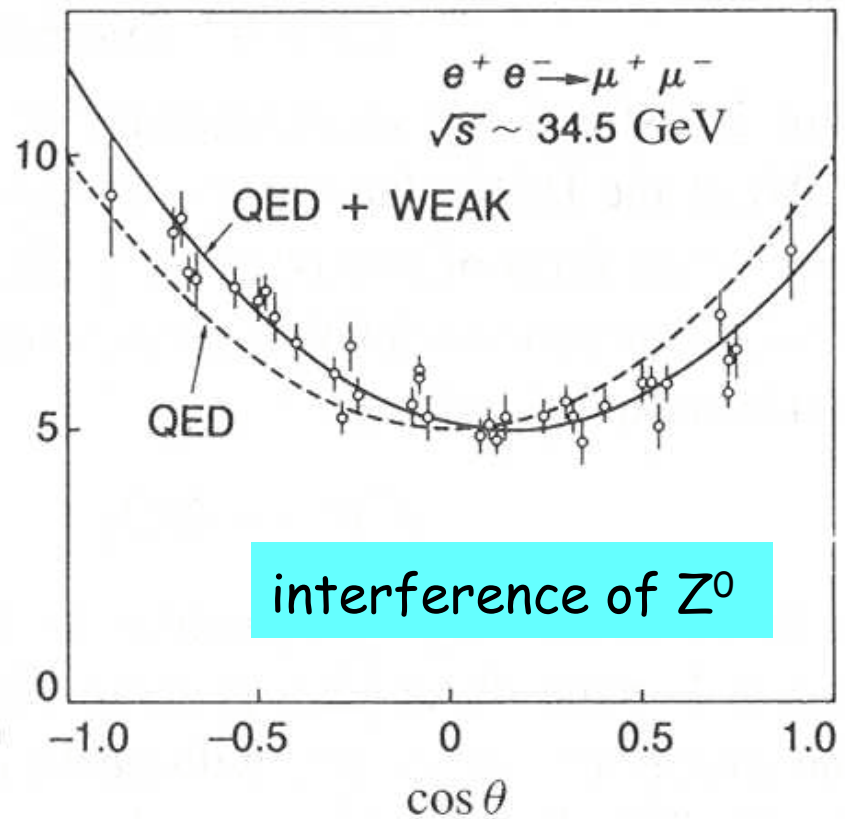
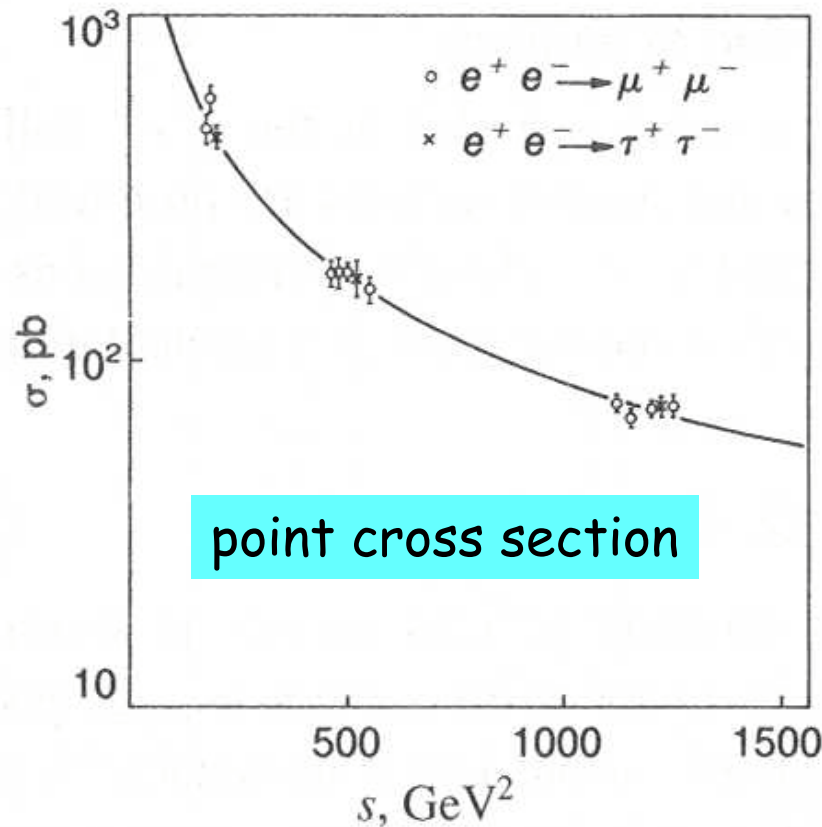
- Amplitude is $d_{mm'}^J(\theta) = d_{1,1}^1(\theta) = (1+\cos \theta)/2$
 - if $RL \rightarrow RL$
- Amplitude is $d_{mm'}^J(\theta) = d_{1,-1}^1(\theta) = (1-\cos \theta)/2$
 - if $RL \rightarrow LR$
- $M^2 \sim [(1+\cos \theta)/2]^2 + [(1-\cos \theta)/2]^2 = (1+\cos^2 \theta)/2$

$$e^+ e^- \rightarrow \mu^+ \mu^-$$

- $\frac{d\sigma}{d\Omega} = \frac{\alpha^2}{4s} (1 + \cos^2 \theta)$
- $\sigma(e^+ e^- \rightarrow \mu^+ \mu^-)$
 $= 4\pi\alpha^2 / 3s$
 $= 87\text{nb} / s(\text{GeV}^2)$
 (point cross section)

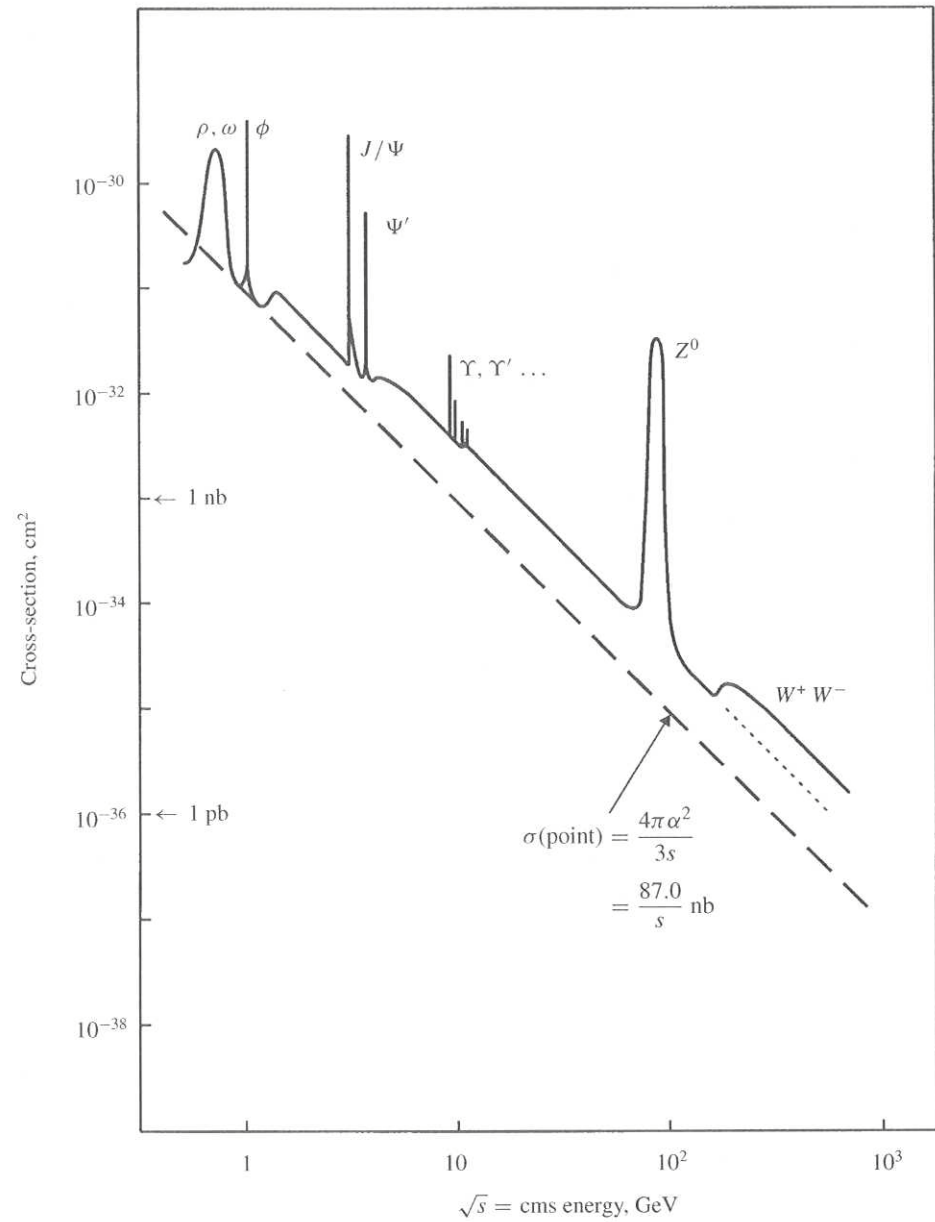
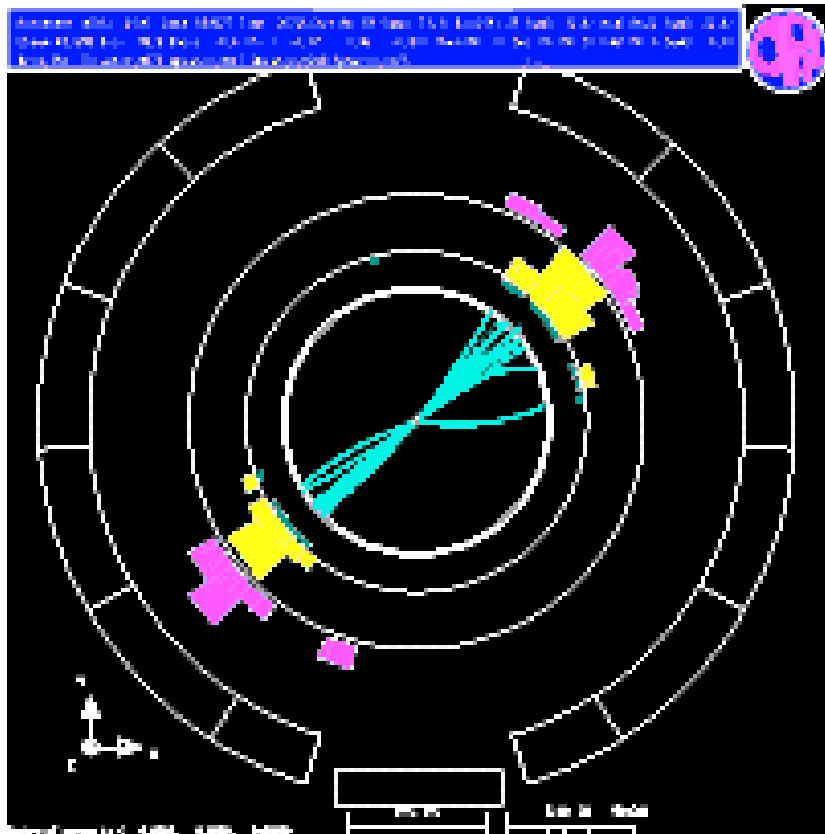


$$e^+ e^- \rightarrow \mu^+ \mu^-$$



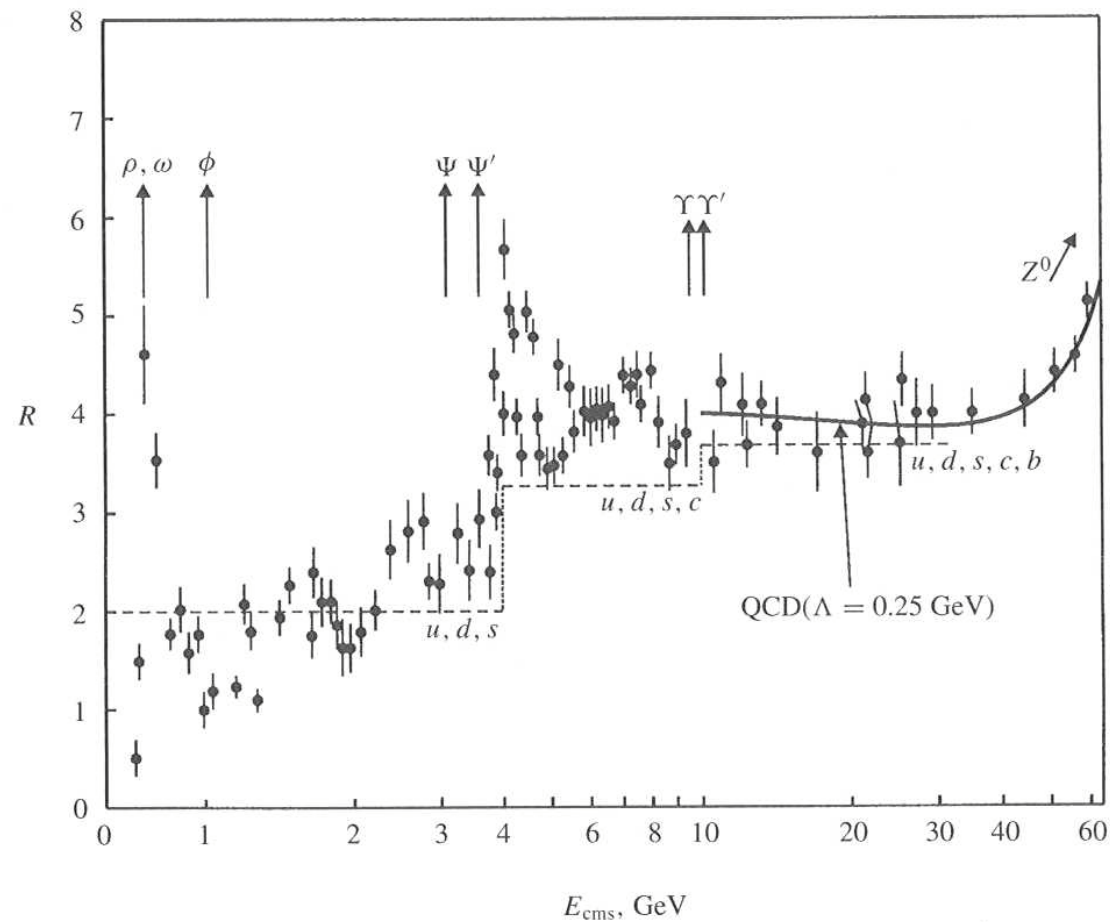
$$f \sim a_{wk} a_{em} / a_{em}^2 \sim Gs / (4\pi\alpha) \sim 10^{-4} s \quad (\text{interference})$$

$e^+ e^-$ annihilation to hadrons



$e^+ e^-$ annihilation to hadrons

- $R = \sigma(e^+ e^- \rightarrow \text{hadrons}) / \sigma(\text{point})$



$e^+ e^-$ annihilation to hadrons

- $R = \sigma(e^+ e^- \rightarrow \text{hadrons}) / \sigma(\text{point})$
 - consider $e^+ e^- \rightarrow \text{hadrons}$ as $e^+ e^- \rightarrow Q\bar{Q}$,
summed over all quarks

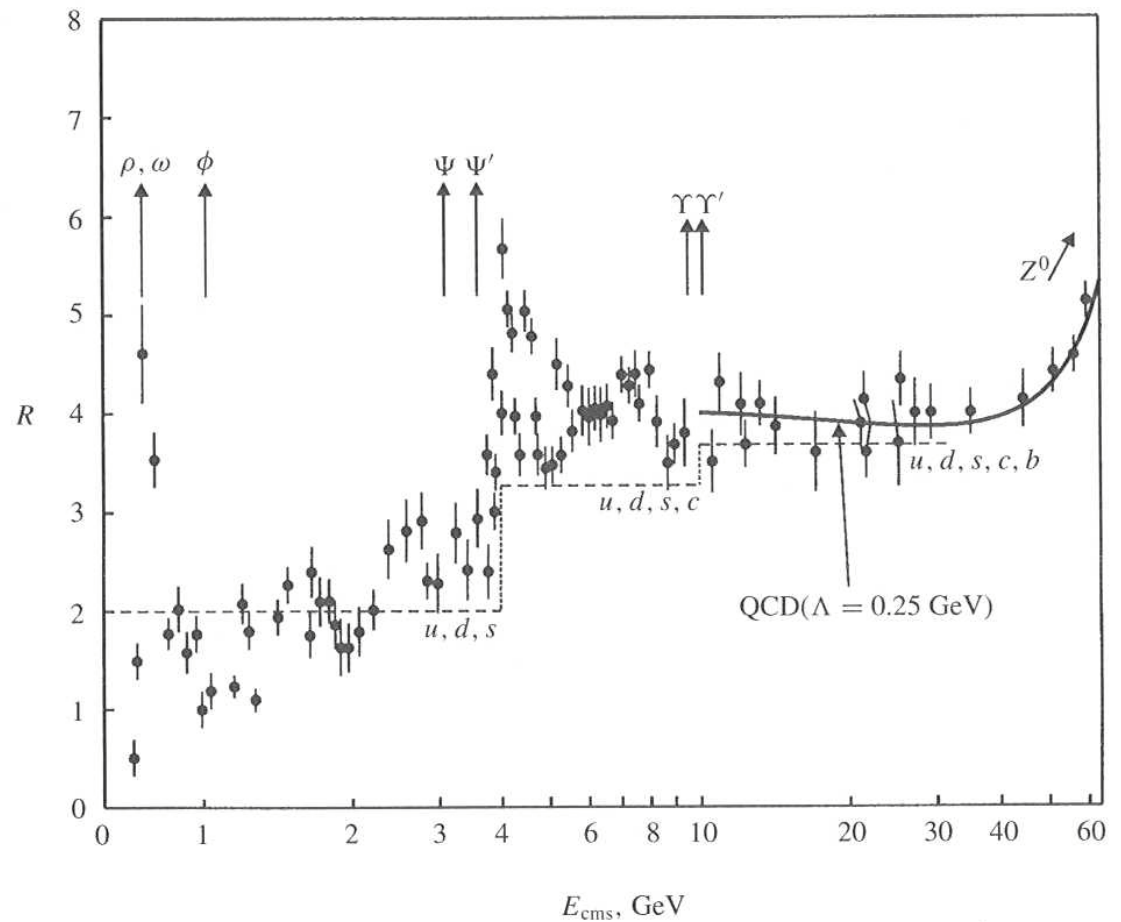
- $R = \sum e_i^2 / e^2$

$$= N_c \left(\underset{d}{(1/3)^2} + \underset{u}{(2/3)^2} + \underset{s}{(1/3)^2} + \underset{c}{(2/3)^2} + \underset{b}{(1/3)^2} + \dots \right)$$
 - $N_c = 3$

- Therefore, R should increase by a well defined value as each flavor threshold is crossed

$e^+ e^-$ annihilation to hadrons

Threshold	R
below c	2
charm	$3 \frac{1}{3}$
bottom	$3 \frac{2}{3}$

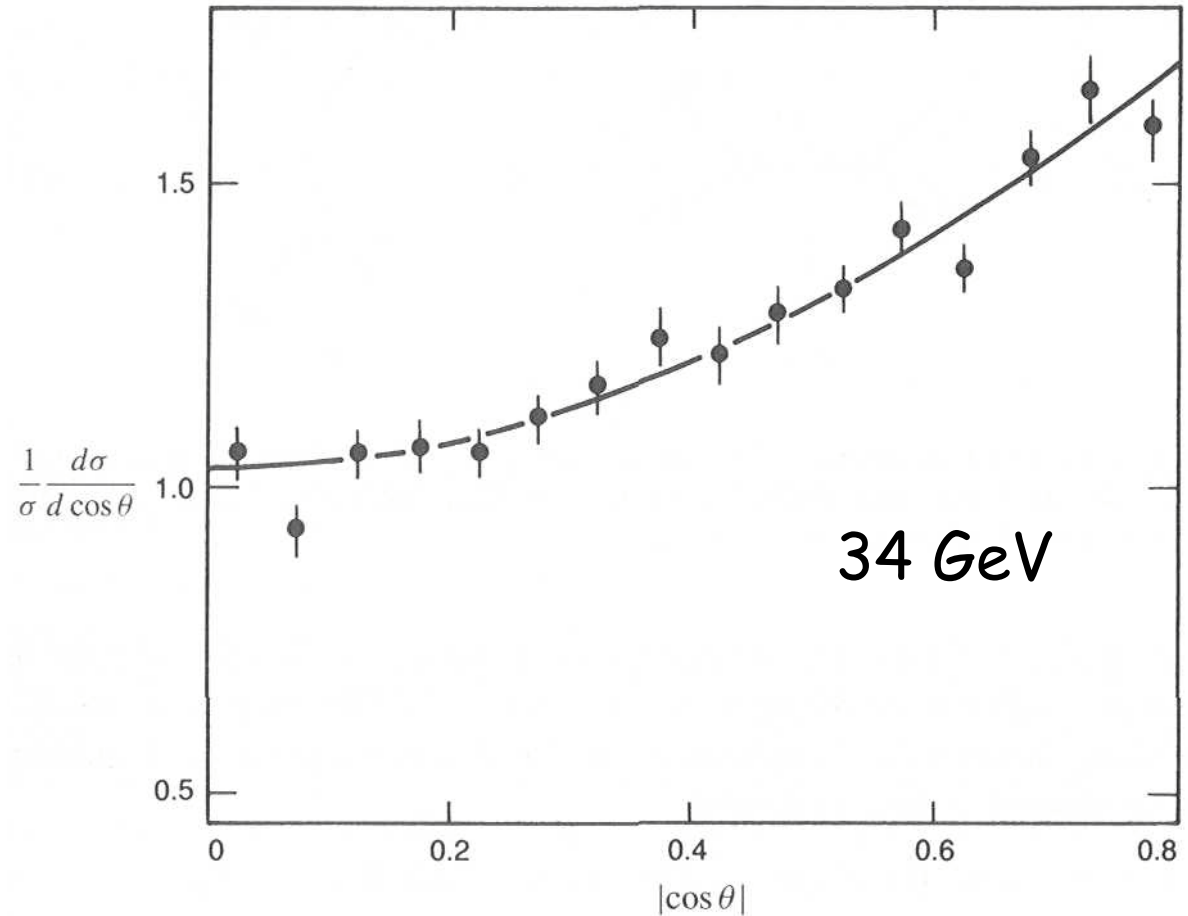


$e^+ e^-$ annihilation to hadrons

$$e^+ e^- \rightarrow Q \bar{Q}$$

If Q is spin 1/2

$$(1 + \cos^2 \theta)$$

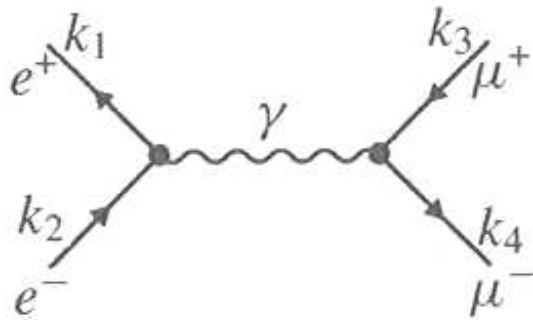


$e^+ e^-$ annihilation to hadrons

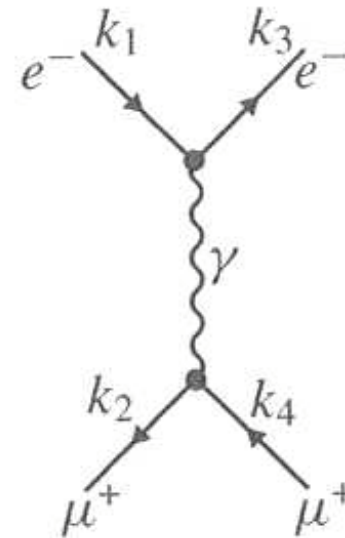
- Constancy of R indicates pointlike constituents
- Angular distributions of hadron jets prove spin $1/2$ partons
- Values of R consistent with partons of quarks expected charge and color quantum number

Electron-muon scattering, $e^- \mu^+ \rightarrow e^- \mu^+$

s channel



t channel



- Mandelstam variables:

$$s = -(k_1 + k_2)^2 = -(k_3 + k_4)^2 = -2k_1k_2 = -2k_3k_4$$

$$t = q^2 = (k_1 - k_3)^2 = (k_2 - k_4)^2 = -2k_1k_3 = -2k_2k_4$$

$$u = (k_2 - k_3)^2 = (k_1 - k_4)^2 = -2k_2k_3 = -2k_1k_4$$

for $m = 0$

$$k = (p_x, p_y, p_z, iE)$$

Electron-muon scattering, $e^- \mu^+ \rightarrow e^- \mu^+$

- Mandelstam variables:

$$s = -(k_1 + k_2)^2 = -(k_3 + k_4)^2 = -2k_1 k_2 = -2k_3 k_4$$

$$\text{suppose } k_1 = (p, iE) \quad k_2 = (-p, iE)$$

$$s = -(p-p)^2 - (2iE)^2 = 4E^2 \quad \text{for } m = 0$$

$$t = q^2 = (k_1 - k_3)^2 = (k_2 - k_4)^2 = -2k_1 k_3 = -2k_2 k_4$$

$$\text{suppose } k_3 = (p \cos \theta, p \sin \theta, 0, iE)$$

$$\begin{aligned} t &= (p - p \cos \theta)^2 + (-p \sin \theta)^2 + (E - E)^2 = 2p^2(1 - \cos \theta) \\ &= 4p^2 \sin^2 \theta / 2 \end{aligned}$$

$$u = (k_2 - k_3)^2 = (k_1 - k_4)^2 = -2k_2 k_3 = -2k_1 k_4 = 4p^2 \cos^2 \theta / 2$$

Electron-muon scattering, $e^- \mu^+ \rightarrow e^- \mu^+$

- Annihilation cross-section ($e^+e^- \rightarrow \mu^+\mu^-$)
- $\frac{d\sigma}{d\Omega} = \frac{\alpha^2}{8p^2} \left(\frac{t^2 + u^2}{s^2} \right) = \frac{\alpha^2}{8p^2} [\sin^4(\theta/2) + \cos^4(\theta/2)]$

$$= \frac{\alpha^2}{4s} [1 + \cos^2 \theta]$$

$$\begin{aligned} 2 \sin^2(\theta/2) &= 1 - \cos \theta \\ 2 \cos^2(\theta/2) &= 1 + \cos \theta \end{aligned}$$

$$s = 4p^2$$

- Now consider the crossed channel, $e^- \mu^+ \rightarrow e^- \mu^+$

Electron-muon scattering, $e^- \mu^+ \rightarrow e^- \mu^+$

crossed channel ($t \leftrightarrow u$)

$$\bullet \quad \frac{d\sigma}{d\Omega} = \frac{\alpha^2}{8p^2} \left(\frac{s^2 + u^2}{t^2} \right)$$

$$= \frac{\alpha^2}{8p^2 \sin^4(\theta/2)} [1 + \cos^4(\theta/2)]$$

So far, all of these expressions are for the center-of-momentum system

What about the laboratory frame?

Electron-muon scattering, $e^- \mu^+ \rightarrow e^- \mu^+$

- Suppose the muon is a target at rest in the laboratory
(Note: this is not practical since the lifetime of the muon is 2.2 microseconds)
- γ is the boost from cms to lab
- p and θ are the projectile parameters in the cms

- In the lab:

$$\begin{aligned} E_\mu &= \gamma (p - \beta p \cos \theta) \quad (\text{scattered muon}) \\ &= \gamma p (1 - \cos \theta) \end{aligned}$$

$$E_e = \gamma (p + \beta p) = 2\gamma p \quad (\text{incident electron})$$

$$\gamma = E_\mu / E_e = (1 - \cos \theta) / 2$$

$$\text{so } \cos^2 \theta / 2 = (1 + \cos \theta) / 2 = 1 - \gamma$$

$$d\Omega = 2\pi d(\cos \theta) = 4\pi d\gamma$$

Electron-muon scattering, $e^- \mu^+ \rightarrow e^- \mu^+$

$$\frac{d\sigma}{d\Omega} = \frac{\alpha^2}{8p^2 \sin^4(\theta/2)} [1 + \cos^4(\theta/2)]$$

$d\Omega = 4\pi dy$
 $\cos^2 \theta/2 = 1-y$
 $q^2 = t = 4p^2 \sin^2 \theta/2$

$$\frac{d\sigma}{4\pi dy} = \frac{\alpha^2}{q^4 / 2p^2} [1 + (1-y)^2]$$

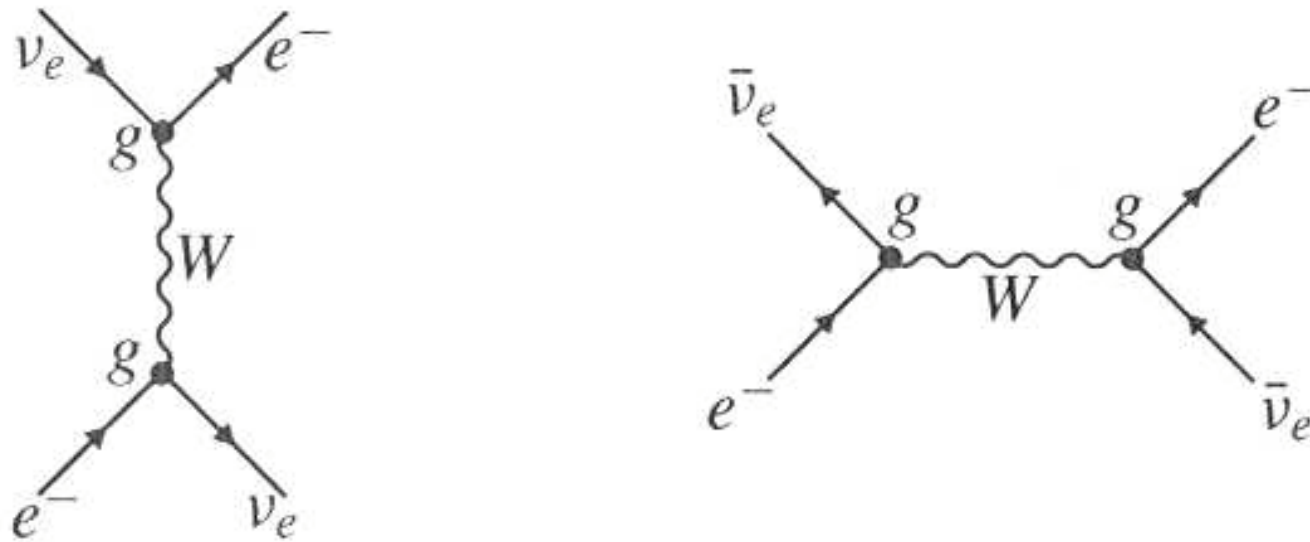
same helicity scattering

$$\frac{d\sigma}{dy} = \frac{2\pi \alpha^2 s}{q^4} [1 + (1-y)^2]$$

opposite helicity scattering

As $y \rightarrow 0$, we recover the Rutherford formula

Neutrino-electron scattering, $\nu_e e^- \rightarrow \nu_e e^-$



Consider only the charged-current reaction

$$\begin{aligned} M(\nu_e e^- \rightarrow \nu_e e^-) &= (g/\sqrt{2})^2 / (q^2 + M_w^2) \\ &= g^2 / 2M_w^2 \quad \text{for } q^2 \ll M_w^2 \end{aligned}$$