

2.3.1. **Quadrupole moments** (7 pts for parts a) - d) plus 3 pts for part e))

- a) Consider a localized charge density as in ch.2 §3.1 and carry the expansion of the potential to $O(1/r^3)$. Show that the potential to that order is given by

$$\varphi(\mathbf{x}) = \frac{1}{r} Q + \frac{1}{r^3} \mathbf{x} \cdot \mathbf{d} + \frac{1}{r^5} \sum_{i,j} x_i x_j Q_{ij} + \dots$$

with Q the total charge and $\mathbf{d}$ the dipole moment, and determine the quadrupole tensor Q_{ij} .

- b) Show that the quadrupole tensor is independent of the choice of the origin provided the total charge and the dipole moment vanish.
- c) Consider a homogeneously charged ellipsoid $(x/a)^2 + (y/b)^2 + (z/c)^2 \leq 1$ and calculate the quadrupole tensor Q_{ij} with respect to the ellipsoid's center. Check to make sure that the result for Q_{ij} is traceless.
- d) Let the charge density be invariant under rotations about the z -axis through multiples of an angle α , with $|\alpha| < \pi$. Show that in this case the quadrupole tensor has the form $\begin{pmatrix} q & 0 & 0 \\ 0 & q & 0 \\ 0 & 0 & -2q \end{pmatrix}$. Make sure your result from part c) conforms with this for the special case $a = b$.
- e) Consider the homogeneously charged ellipsoid from Part c), and calculate the quadrupole moments Q_{2m} as defined in ch.2 §3.5.

(7 points)

$$\begin{aligned}
 2.1.1. a) \quad \frac{1}{|\vec{x}-\vec{y}|} &= \frac{1}{r} \left(1 - 2 \frac{\vec{x} \cdot \vec{y}}{r^2} + \frac{y^2}{r^2} \right)^{-1/2} = \frac{1}{r} \left(1 + \frac{\vec{x} \cdot \vec{y}}{r^2} - \frac{1}{2} \frac{y^2}{r^2} + \frac{3}{2} \frac{(\vec{x} \cdot \vec{y})^2}{r^4} + \dots \right) \\
 &= \frac{1}{r} \left[1 + \frac{\vec{x} \cdot \vec{y}}{r^2} + \frac{3}{2} x_i x_j y_i y_j \frac{1}{r^4} - \frac{1}{2} y^2 \delta_{ij} x_i x_j \frac{1}{r^4} + \dots \right] \\
 &= \frac{1}{r} \left[1 + \frac{\vec{x} \cdot \vec{y}}{r^2} + \frac{1}{2} x_i x_j (3 y_i y_j - \delta_{ij} y^2) \frac{1}{r^4} + \dots \right]
 \end{aligned}$$

$$\begin{aligned}
 \Rightarrow \underline{\underline{\varphi(\vec{x})}} &= \int d\vec{y} \frac{\rho(\vec{y})}{|\vec{x}-\vec{y}|} = \frac{1}{r} \int d\vec{y} \rho(\vec{y}) + \frac{1}{r^2} \vec{x} \cdot \int d\vec{y} \vec{y} \rho(\vec{y}) \\
 &\quad + \frac{1}{2} \frac{1}{r^2} x_i x_j \int d\vec{y} (3 y_i y_j - \delta_{ij} y^2) \rho(\vec{y}) + \dots \\
 &= \underline{\underline{\frac{1}{r} Q + \frac{1}{r^2} \vec{x} \cdot \vec{d} + \frac{1}{r^2} \sum_{ij} x_i x_j Q_{ij} + O(1/r^4)}}
 \end{aligned}$$

also $Q = \int d\vec{y} \rho(\vec{y})$ monopole moment

$\vec{d} = \int d\vec{y} \vec{y} \rho(\vec{y})$ dipole moment

$\underline{\underline{Q_{ij} = \frac{1}{2} \int d\vec{y} (3 y_i y_j - \delta_{ij} y^2) \rho(\vec{y})}}$ quadrupole moment

b) $\rho'(\vec{y}) = \rho(\vec{y} - \vec{a})$

$$\begin{aligned}
 \Rightarrow \underline{\underline{Q'_{ij}}} &= \frac{1}{2} \int d\vec{y} (3 y_i y_j - \delta_{ij} y^2) \rho'(\vec{y}) \\
 &= \frac{1}{2} \int d\vec{y} [3 (y_i + a_i)(y_j + a_j) - \delta_{ij} (\vec{y} + \vec{a})^2] \rho(\vec{y}) \\
 &= Q_{ij} + \frac{1}{2} \int d\vec{y} [3 a_i y_j + 3 a_j y_i + 3 a_i a_j - \delta_{ij} (2 \vec{a} \cdot \vec{y} + a^2)] \rho(\vec{y}) \\
 &= Q_{ij} + \frac{3}{2} a_i d_j + \frac{3}{2} a_j d_i + \frac{3}{2} a_i a_j Q - \delta_{ij} \vec{a} \cdot \vec{d} - \delta_{ij} \frac{1}{2} a^2 Q \\
 &= \underline{\underline{Q_{ij}}} \quad \text{if} \quad \underline{\underline{\vec{d} = a = 0}}
 \end{aligned}$$

c) ellipsoid. $x^2/a^2 + y^2/b^2 + z^2/c^2 \leq 1$

$$\rightarrow Q_{ij} = \frac{1}{2} \int d\vec{x} (\partial x_i \partial x_j - \delta_{ij} \nabla^2) \Theta(x^2/a^2 + y^2/b^2 + z^2/c^2 \leq 1) \int$$

where $\vec{x} = (x, y, z)$ and $\int = \text{volume element}$

by symmetry $\rightarrow \underline{\underline{\delta_{ij} = 0 \text{ unless } i=j}}$

$$\underline{\underline{Q_{11}}} = \frac{8}{2} \int dx dy dz (x^2 - y^2 - z^2) \Theta(x^2/a^2 + y^2/b^2 + z^2/c^2 \leq 1)$$

$$= \frac{1}{2} \int abc \int dx dy dz (2a^2 x^2 - b^2 y^2 - c^2 z^2) \Theta(x^2/a^2 + y^2/b^2 + z^2/c^2 \leq 1)$$

$$= \frac{1}{2} \int abc \left[2a^2 \int_0^1 dr r^2 \int_{-1}^1 dz \int_0^{2\pi} d\phi r^2 \sin^2 \theta \right. \\ \left. - b^2 \int_0^1 dr r^2 \int_{-1}^1 dz \int_0^{2\pi} d\phi r^2 \sin^2 \theta \cos^2 \theta \right. \\ \left. - c^2 \int_0^1 dr r^2 \int_{-1}^1 dz \int_0^{2\pi} d\phi r^2 \cos^2 \theta \right]$$

$$= \frac{1}{2} \int abc \left[2a^2 \int_0^1 dr r^2 (2 - \frac{2}{3}) - b^2 \int_0^1 dr r^2 (2 - \frac{2}{3}) - c^2 \int_0^1 dr r^2 (2 - \frac{2}{3}) \right]$$

$$= \frac{1}{2} \int abc \frac{4\pi}{3} \left[\frac{8}{3} a^2 - \frac{4}{3} b^2 - \frac{4}{3} c^2 \right]$$

$$= \frac{4\pi}{3} abc \int_0^1 \frac{1}{10} (2a^2 - b^2 - c^2)$$

$$= \underline{\underline{Q \frac{1}{10} (2a^2 - b^2 - c^2)}}$$

with $\underline{\underline{Q = \frac{4\pi}{3} \int abc = \text{total volume}}}$

$$\underline{\underline{Q_{22}}} = \underline{\underline{Q \frac{1}{10} (2b^2 - a^2 - c^2)}}$$

by symmetry

$$\underline{\underline{Q_{33}}} = \underline{\underline{Q \frac{1}{10} (2c^2 - a^2 - b^2)}}$$

by symmetry

check. $\underline{\underline{Q_{11} + Q_{22} + Q_{33} = 0}}$ ✓

d) As a real symmetric tensor, Q_{ij} can always be diagonalized
 $\rightarrow$ The most general form of Q_{ij} in its principal axes system is

$$Q_{ij} = \begin{pmatrix} q_+ + q_- & 0 & 0 \\ 0 & q_+ - q_- & 0 \\ 0 & 0 & -2q_+ \end{pmatrix}$$

where

$$q_- = \frac{1}{2} (Q_{11} - Q_{22}) = \frac{1}{2} \int d\vec{x} \, f(\vec{x}) [2x^2 - y^2 - z^2 - (y^2 + x^2 + z^2)]$$

$$= \frac{3}{2} \int d\vec{x} \, f(\vec{x}) (x^2 - y^2)$$

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cylindrical coordinates:

$$x = r \cos \varphi \\ y = r \sin \varphi$$

$$x^2 - y^2 = r^2 (\cos^2 \varphi - \sin^2 \varphi) = r^2 \cos 2\varphi$$

$$\rightarrow q_- = \frac{3}{2} \int_0^{2\pi} d\varphi \int_0^\infty dr r \int_{-\infty}^\infty dz \, f(r, \varphi, z) r^2 \cos 2\varphi$$

$$\text{Now let } f(r, \varphi, z) = f(r, \varphi + \alpha, z)$$

$$\begin{aligned} \rightarrow q_- &= \frac{3}{2} \int_0^{2\pi} d\varphi \int_0^\infty dr r^2 \int_{-\infty}^\infty dz \, f(r, \varphi + \alpha, z) \cos 2\varphi \\ &= \frac{3}{2} \int_0^{2\pi} d\varphi \int_0^\infty dr r^2 \int_{-\infty}^\infty dz \, f(r, \varphi, z) \cos 2(\varphi + \alpha) \\ &= \frac{3}{2} \int_0^{2\pi} d\varphi \int_0^\infty dr r^2 \int_{-\infty}^\infty dz \, f(r, \varphi, z) (\cos 2\varphi \cos 2\alpha - \sin 2\varphi \sin 2\alpha) \end{aligned}$$

$$\boxed{r^2 \sin 2\varphi = r^2 \sin \varphi \cos \varphi = xy \rightarrow \text{the second term is } \propto Q_{12} = 0}$$

$$= \cos 2\alpha \cdot q_- \rightarrow \underline{q_- = 0} \text{ when } \alpha \neq 0$$

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part c) with $a=b \rightarrow Q_{11} = Q_{22} = \frac{Q}{10} (a^2 - c^2) \checkmark$

$$p = 2, 1, 1, -4$$

$$e.) \quad \underline{Q_{20}} = \sqrt{\frac{4\pi}{5}} \int_0^\infty dr r^4 \int dR g(r, R) \sqrt{\frac{5}{4\pi}} \frac{1}{2} (1 - \gamma^2) \\ = \frac{1}{2} \int d\vec{x} g(\vec{x}) (\gamma^2 - r^2) = \underline{Q_{22}}$$

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$$\underline{Q_{2, \pm 2}} \propto \underbrace{\int dR g(r, R)}_{\text{even}} \underbrace{P_2^{\pm 2}(\gamma)}_{\text{odd}} = 0 \quad \text{by symmetry}$$

tot. of γ

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$$\underline{Q_{22}} = \sqrt{\frac{4\pi}{5}} \int_0^\infty dr r^4 \int dR g(r, R) \sqrt{\frac{5}{4\pi}} \frac{1}{4!} e^{2i\varphi} \gamma^2 (1 - \gamma^2) \\ = \frac{1}{124} \int d\vec{x} g(\vec{x}) r^2 (1 - \gamma^2) (\cos 2\varphi + i \sin 2\varphi) \\ = \frac{1}{124} \int d\vec{x} g(\vec{x}) r^2 \underbrace{\cos^2 \varphi}_{\text{even}} (\cos^2 \varphi - \sin^2 \varphi) + \frac{1}{124} \int d\vec{x} g(\vec{x}) r^2 \underbrace{\sin^2 \varphi}_{\text{even}} (\cos^2 \varphi - \sin^2 \varphi) \\ = \frac{1}{124} \int d\vec{x} g(\vec{x}) r^2 (\underbrace{\cos^2 \varphi - \sin^2 \varphi}_{\text{tot. of } \varphi})$$

$$\left. \begin{array}{l} x = r \cos \varphi \\ y = r \sin \varphi \\ z = r \cos \varphi \end{array} \right\} \rightarrow x^2 - y^2 = r^2 \cos^2 \varphi - r^2 \sin^2 \varphi$$

$$= \frac{1}{216} \int d\vec{x} g(\vec{x}) (x^2 - y^2) \\ = \frac{1}{216} \int d\vec{x} g(\vec{x}) [(2x^2 - y^2 - z^2) - (y^2 - x^2 - z^2)] \frac{1}{2} \\ = \frac{1}{16} (\underline{Q_{22}} - \underline{Q_{22}})$$

$$\underline{Q_{2,-2}} = \sqrt{\frac{4\pi}{5}} \int_0^\infty dr r^4 \int dR g(r, R) \sqrt{\frac{5}{4\pi}} \frac{1}{4!} e^{-2i\varphi} \gamma^2 (1 - \gamma^2) = \underline{Q_{22}}$$

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