

Quantum field theory = states of a system
are sections of a bundle
We're interested in a particular 3-d theory
(so underlying spaces are 3-manifolds) constructed
from a group G , ~~and~~ (cpt) and a G -rep V .
Part of data is a principal G -bundle
 P on 3-mfld M , and a section ϕ of
associated bundle $(P \times V)/G \rightarrow M$.

Coulomb branch = Spec of algebra of "local
operators" = Spec of Hilbert space of S^2

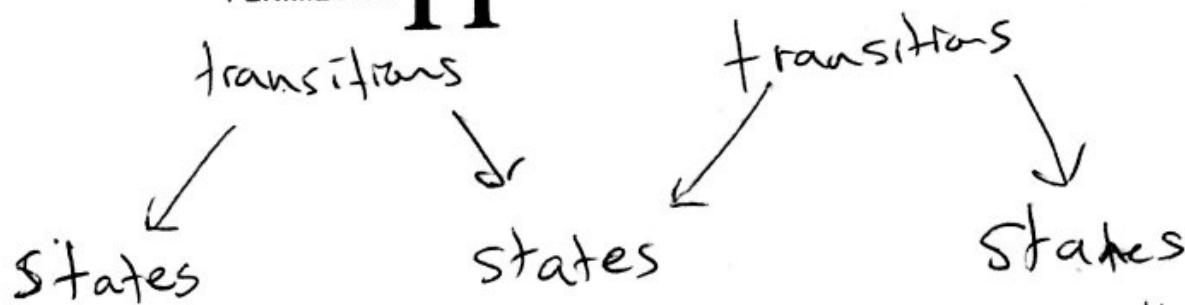
This is a commutative algebra because at
 considered as a cobordism.

expected value of operator = "integral" over "states"
 $\Rightarrow$ Hilbert space of surface = $H_*(\text{ways of translation}$
invariant states on $C \times \mathbb{R}$).

$\Rightarrow$ local operators = H_* (ways of transitioning
from one state to another, locally).

This has a product structure via diagram





Since this is local we can think about $\mathbb{R}^3 = \mathbb{R}^2 \times \mathbb{R}$. We look at the way that (space time) we can change the principal bundle and section at the origin from one at $t = -\epsilon$ and $t = \epsilon$, which is constant on $\mathbb{R}^2 \setminus \{(0,0)\}$.

We ~~can~~ consider a formal disk P around $(0,0)$. Let $D^* = D \setminus \{0,0\}$. We have a bundle and section (P_1, ϕ_1) at $t = -\epsilon$ and one (P_2, ϕ_2) at $t = \epsilon$. These are isomorphic on D^* .

Let us choose a trivialization of this bundle on D^* (which carries from $t = -\epsilon$ to $t = \epsilon$).

The only redundancy in this description is the choice of trivialization, which are a torsor over $G(t) = \text{Maps}(D^* \rightarrow G)$. The choice of a principal bundle on D and a trivialization on D^* is parametrized by the affine Grassmannian $G(t)$ acts transitively on the ~~parametrizations~~ trivializations.



The ~~obvious trivialization~~ is preserved by ~~G(t)~~. The re-trivializations of the standard trivialization that extend our D are given by $G(t)$. Adding a choice of ϕ is an ~~choice~~ element of the associated bundle $(G(t) \times V(t)) / G(t) = \mathcal{X}_V$ we need to have

ϕ_1 and ϕ_2 agree on D^* ; thinking of these as elements of $V(t)$ using the trivialization, B the map $\mathcal{X}_V \rightarrow V(t)$

So, the choice of (P_i, ϕ_i) with $\phi_1 \neq \phi_2$ and the trivialization B given by the trivialization $\mathcal{X}_V \times_{V(t)} \mathcal{X}_V$; forgetting the trivialization B modding out by $G(t)$.

So we need to consider $H_*^{G(t)}(\mathcal{X}_V \times_{V(t)} \mathcal{X}_V)$
This is the algebra of local operators

The map $\mathcal{X}_V \times_{V(t)} \mathcal{X}_V \rightarrow \mathcal{X}_V \rightarrow \mathcal{B}$ is $G(t)$ -equivariant. Let $v \in \mathcal{X}_V$ be the fiber over the identity coset $v \in \mathcal{X}_V = \{(g(t), \phi) \mid g(t) \cdot \phi \in V(t)\}$

$$H_*^{G(t)}(\mathcal{X}_V \times_{V(t)} \mathcal{X}_V) \cong H_*^{G(t)}(v \mathcal{X}_V)$$



What does this homology actually mean?
 Don't think too hard about it, but here's the rough idea: $\mathcal{H}_v$ is an infinite dimensional vector bundle on an infinite dimensional spaces.

We want homology classes "finite along base, infinite along fiber." Actually making this precise is quite annoying: $\mathcal{H}_v$ is an inverse limit of f.d. v.s. $V[t]/t^N$ as $N \rightarrow \infty$. The associated bundles of $V[t]/t^N$ as $N \rightarrow \infty$. Every question we consider is insensitive to high enough powers of N . Thus, we can think about the Borel Moore homology in the usual sense of these truncations and then identify them using the pull back isomorphism for N . But don't let this scare you. We just need to use a few simple facts:

1) a finite dimensional subvariety in G_0 , ~~with~~ with a coherent rank sub bundle induces a homology class.

2) ~~pushing forward by a smooth inclusion and then pulling back has the effect of multiplying by the Euler class of the normal bundle.~~

3) a clean intersection of homology classes (one that is locally diff. to v.s.) results in the usual intersection fines the Euler class of the excess.



How does this work for a torus?
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$$1 \mapsto T_1[[t]] \rightarrow T[[t]] \rightarrow T \rightarrow 1$$

The group $T_1[[t]]$ is contractible.

$$1 \rightarrow T[[t]] \rightarrow T((t)) \rightarrow X_*(T) \rightarrow 1$$

$$D((t)) = \begin{bmatrix} a_1 & & 0 \\ & \ddots & \\ 0 & & a_n \end{bmatrix}$$

$$a_i(t) = b_{k_i}^{(i)} t^{k_i} + b_{k_i+1}^{(i)} t^{k_i+1} + \dots$$

$b_{k_i}^{(i)} \neq 0$

$$D[[t]] = \begin{bmatrix} a_1 & & 0 \\ & \ddots & \\ 0 & & a_n \end{bmatrix} \quad k_i = 0$$

$$X_*(T) \begin{bmatrix} k_1 & & 0 \\ & \ddots & \\ 0 & & k_n \end{bmatrix} \leftarrow \begin{bmatrix} a_1 & & 0 \\ & \ddots & \\ 0 & & a_n \end{bmatrix}$$

$$D_1[[t]] \begin{bmatrix} a_1 & & 0 \\ & \ddots & \\ 0 & & a_n \end{bmatrix} \quad k_i = 0, \quad b_0^{(i)} = 1$$

Thus $T((t)) / T[[t]] \cong X_*(T)$
 $[t^v]$ for $\gamma \in X_*(T)$



Let $V_0 = V \times \{x\}$. We have

$$\begin{aligned} \mathcal{X}_V &= \bigsqcup_{\lambda \in X_*(T)} \{\lambda^\vee\} \times \lambda^\vee \cdot V_0 \\ \vee \mathcal{X}_V &= \bigsqcup_{\lambda \in X_*(T)} \{\lambda^\vee\} \times (V_0 \cap \lambda^\vee \cdot V_0) \end{aligned}$$

As an ~~abelian~~ ^{abelian} group, there's nothing interesting about the homology. Let $\gamma_\lambda = [\{\lambda^\vee\} \times (V_0 \cap \lambda^\vee \cdot V_0)]$

$$H_*^{G_0}(\vee \mathcal{X}_V) = \oplus H^*(BT) \cdot \gamma_\lambda$$

The only interesting thing is the multiplication. Left and right multiplication by $H^*(BT)$ corresponds to the tautological bundles on the left and right copies of $\mathcal{X}_V$. Transferring to $\vee \mathcal{X}_V$, right multiplication corresponds to the Euler classes of tautological bundles, and left to the trivial bundle with the character giving the equivariant structure.

In this case, ~~the~~ the tautological bundles are trivial, so the left and right actions are the same.



So, the only interestingly product is $r_2 \gamma$
 we pull back their classes to
 $X_v \times X_v \times X_v$, and then intersect them

$$p_{12}^* r_2 = \left[\begin{array}{c} (x_1 + \lambda, x_2, x_3) \\ (v_1(t), v_2(t), v_3(t)) \end{array} \right]$$

$$p_{23}^* r_2 = \left[\begin{array}{c} (\lambda_1, \lambda_2 + 2, \lambda_3) \\ (w_1(t), w_2(t), w_3(t)) \end{array} \right]$$

The naive intersection is

$$\left[\begin{array}{c} (x + \lambda + v, x + v, \lambda) \\ (v(t), v(t), v(t)) \end{array} \right]$$

$$w/ \begin{array}{c} t^{x+\lambda+v} v_0 \\ t^{x+v} v_0 \\ t^x v_0 \end{array}$$

However, we have to account for the excess bundle, given by

$$t^{x+\lambda+v} v_0 + t^{x+v} v_0 + t^x v_0$$

$$(t^{x+\lambda+v} v_0 + t^{x+v} v_0 + t^x v_0)$$

If v is of weight μ , then $t^n v \in \mathcal{V}_0$ if $n \geq \langle x, \mu \rangle$. Thus $t^n v$ contributes to the excess if $n \geq \langle x + v, \mu \rangle$ and $n < \langle x, \mu \rangle$ and $n < \langle x + \lambda + v, \mu \rangle$, and we get a factor of μ for each such n .



Note, this requires $\langle \lambda, \mu \rangle < 0$ and $\langle \lambda, \mu \rangle > 0$.

In addition, we must push forward. By the map forgetting the middle element this sends us to the class of $\begin{pmatrix} t^{x+\lambda+\mu}, x \\ v(t), v(t) \end{pmatrix}$. $v(t) \in t^{x+\lambda} v_0 \cap t^x v_0$

This is almost $r_{\lambda+\mu}$, but we have the

$$t^{x+\lambda+\mu} v_0 \cap t^x v_0 / t^{x+\lambda+\mu} v_0 \cap t^{x+\mu} v_0 \cap t^x v_0$$

Now, we need for $t^x v$ to contribute there, we must have $\langle \lambda+\mu, n \rangle$ and $\langle \lambda, n \rangle$ and $\langle \lambda+\mu+\lambda, n \rangle$. This requires $\langle \lambda, \mu \rangle > 0$ and $\langle \lambda, \mu \rangle < 0$.

Thus, we have $r_\lambda r_\mu = r_{\lambda+\mu} \cdot \prod e_i$ where $d(k, l) = \begin{cases} 0 & k, l \text{ have same sign} \\ \min(|k|, |l|) & k, l \text{ have different signs} \end{cases}$



This is the multiplication rule in the classical limit of the hyperbolic enveloping algebra. In order to get the algebra itself, we ~~let~~ go back to the physics. There, we quantize using an Ω -background corresponding to rotation of $\mathbb{R}^2$ in $\mathbb{R}^2 \times \mathbb{R}^1$.

on the affine Grassmannian, this becomes the loop rotation action of $\mathbb{G}^*$.

We have to consider

$$H_*^{\mathbb{G}(K) \times \mathbb{G}^*}(\mathcal{X}_v \times_{v_0} \mathcal{X}_v) \cong H_*^{\mathbb{G}(0) \times \mathbb{G}^*}(\nu \mathcal{X}_v)$$

and we have to remember the loop action on everything. For example

$\mu \cdot r_\lambda$ = tautological bundle for $\mathbb{G}_m$ on $(X+1, X)$ pulled back from the second factor.
Loop $\mathbb{G}^*$ acts w/ weight $\langle X+1, \mu \rangle$

$r_\lambda \cdot \mu =$ first factor
Loop $\mathbb{G}^*$ acts w/ weight $\langle X, \mu \rangle$

Letting h be equivariant parameter for loop $\mathbb{G}^*$



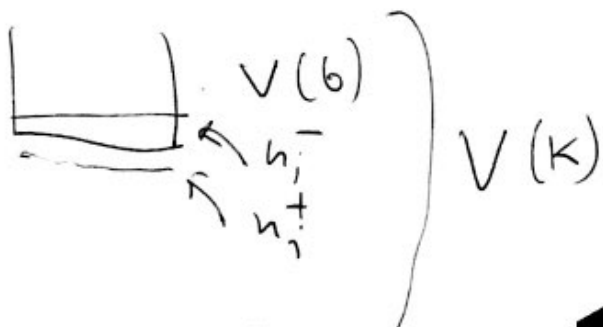
~~This Theorem~~ The convolution algebra $H^*(G) \otimes H^*(G)$ is the classical

This result can be extended in several ways. We can consider equivariance for different groups. The full diagonal matrices H act on V , so we can take $H^*(H \ltimes G) \otimes H^*(G)$ as well.

Secondly, the disk has a $\mathbb{C}^*$ -action by rotation ($\mathbb{R}^2 \cong \mathbb{C}$), which induces an action on $\mathcal{H}_V, \mathcal{H}_V^*$, etc. This ~~is~~ does not commute with $G(K)$; instead, we get an action of the semi-direct product $G(K) \ltimes \mathbb{C}^*$ induced by $s \cdot t^a = s^a t^a$.

I'm going to identify $H^*(B(H \ltimes G)) \cong H^*(B(H \ltimes G) \ltimes \mathbb{C}^*) \cong H^*(B(H \ltimes G) \times \mathbb{C}^*)$

with $D_0[K]$. h_i^- corresponds to the weight of the i th coordinate in $\mathbb{C}^n$ w/ $\mathbb{C}^*$ acting trivially and h_i^+ to the same coordinate w/ $\mathbb{C}^*$ acting w/ weight -1 .



If we turn on the loop $\mathbb{Q}^4$, this is what physicists call an Ω -background and we have to be careful about loop weights, when computing Euler classes. To get right mult, take $\chi = 0$; to get left, $\chi = -\lambda - \nu$.

Consider $\chi = 0$. Then we get

$$r_\lambda r_\nu = r_\lambda r_\nu \cdot \text{Eu}(t^{\lambda+\nu} v_0 + t^\nu v_0 + v_0 / (t^{\lambda+\nu} v_0 \cap t^\nu v_0 + t^\nu v_0 \cap v_0)) \\ \cdot \text{Eu}(t^{\lambda+\nu} v_0 \cap v_0 / t^{\lambda+\nu} v_0 \cap t^\nu v_0 \cap v_0)$$

If $\langle \lambda + \nu, u \rangle \leq 0$, then $t^{\lambda+\nu} v_0 \cap v_0 = v_0$, and a vector of weight u contributes

If $\langle \lambda, u \rangle$ and $\langle \nu, u \rangle$ have same sign, no contribution. If they have opposite sign

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	$\langle \lambda + \nu, \mu \rangle \leq 0$	$\langle \lambda + \nu, \mu \rangle \geq 0$
$\langle \nu, \mu \rangle \geq 0$	$\nu, \dots, t^{\langle \nu, \mu \rangle - 1} \nu$	$t^{\langle \lambda + \nu, \mu \rangle} \nu, \dots, t^{\langle \nu, \mu \rangle - 1} \nu$
$\langle \nu, \mu \rangle < 0$	$t^{\langle \nu, \mu \rangle} \nu, \dots, t^{\langle \lambda + \nu, \mu \rangle + 1} \nu$	$t^{\langle \nu, \mu \rangle}, \dots, t^{-1} \nu$
	$[h_i]^{-\langle \nu, \mu \rangle}$	$[h_i]^{+\langle \nu, \mu \rangle}$ shifted by $\lambda + \nu$
	$[h_i]^{+\langle \lambda, \mu \rangle}$ shifted by $\lambda + \nu$	$[h_i]^{-\langle \nu, \mu \rangle}$

Thus, if we want to be cleaner,
we can move the terms where
 $|\langle \lambda, \mu \rangle| < |\langle \nu, \mu \rangle|$ to the other
side, and get Nick's formula

$$r_\lambda r_\nu = \prod_{\substack{|\langle \lambda, e_i \rangle| < |\langle \nu, e_i \rangle| \\ \langle \lambda, e_i \rangle \langle \nu, e_i \rangle < 0}} [h_i]^{+\langle \lambda, \mu \rangle} r_{\lambda + \nu} [h_i]^{-\langle \nu, \mu \rangle}$$



Thus, we see $H(G) \cong G^*(\vee \mathbb{Z} \vee)$ with $h=1$
 Then H_x is the hyperbolic enveloping algebra
 Note an important subtlety here: my
 r_1 corresponds to a cocharacter
 of T ; Nick's m^2 corresponded to a
 character. Thus, my T is T^V in
 Nick's notation.

(Nick) $1 \rightarrow K \rightarrow (\mathbb{C}^*)^n \rightarrow T \rightarrow 1$

(Me) $1 \rightarrow T^V \rightarrow (\mathbb{C}^*)^n \rightarrow K^V \rightarrow 1$

Switching between $K \subset (\mathbb{C}^*)^n$ and
 $(\mathbb{C}^*)^n / K$ is called Gale duality
 It'll be important for us moving forward

